

## ABINIT School

# Analysis of atomic distortions

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# Born-Oppenheimer energy

**Solid : interacting electrons and nuclei**

- General problem :

$$H(\mathbf{r}, \mathbf{R}) \Phi(\mathbf{r}, \mathbf{R}) = E \Phi(\mathbf{r}, \mathbf{R})$$

↓

$$H(\mathbf{r}, \mathbf{R}) = T_i(\mathbf{R}) + U_{ii}(\mathbf{R}) + T_e(\mathbf{r}) + U_{ee}(\mathbf{r}) + U_{ie}(\mathbf{r}, \mathbf{R})$$

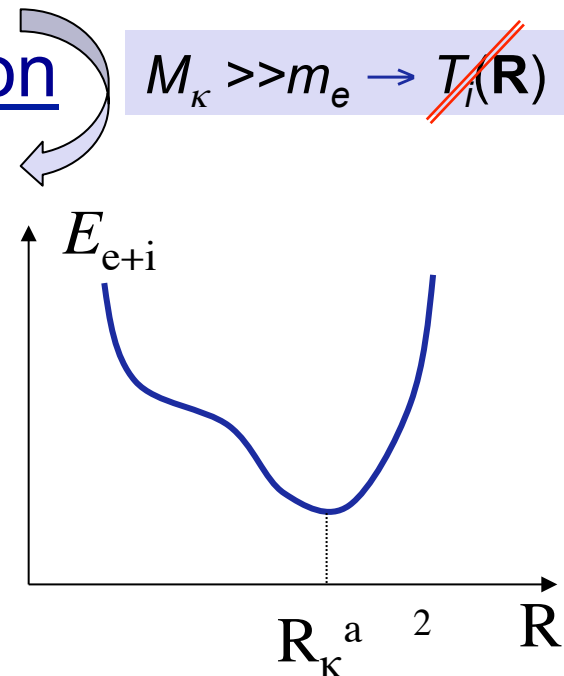
- Born-Oppenheimer approximation

$$M_\kappa \gg m_e \rightarrow \cancel{T_i(\mathbf{R})}$$

$$H_{e+i}(\mathbf{r}, \mathbf{R}) = T_e(\mathbf{r}) + U_{ee}(\mathbf{r}) + U_{ie}(\mathbf{r}, \mathbf{R}) + U_{ii}(\mathbf{R})$$

$$H_{e+i}(\mathbf{r}, \mathbf{R}) \varphi(\mathbf{r}, \mathbf{R}) = E_{e+i}(\mathbf{R}) \varphi(\mathbf{r}, \mathbf{R})$$

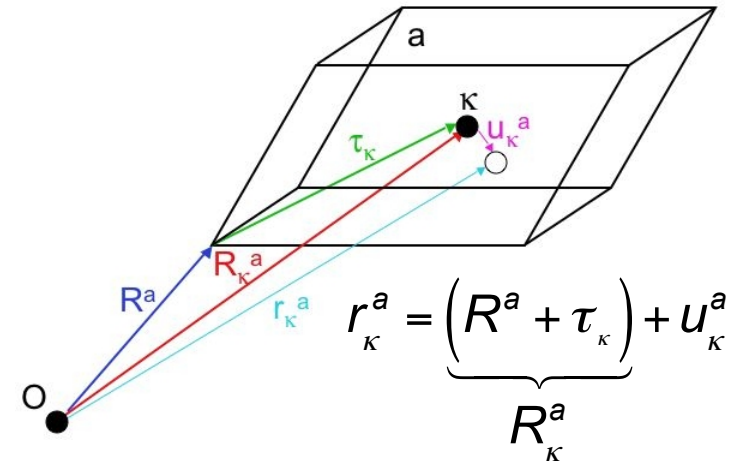
ABINIT total energy  $E_{\text{tot}}(\mathbf{R})$   
corresponds to the Born-  
Oppenheimer energy  $E_{e+i}(\mathbf{R})$



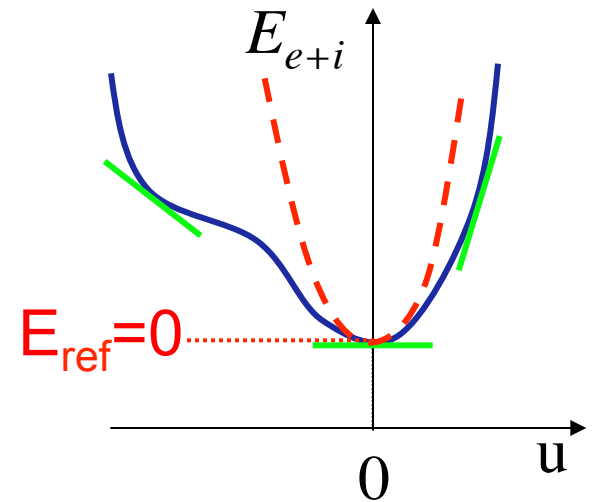
# Background

# Phonons: harmonic approximation

$$E_{e+i}(r_{\kappa}^a) = E_{e+i}(R_{\kappa}^a) - \sum_{\kappa\alpha a} \underbrace{F_{\kappa\alpha}^a}_{\frac{\partial E}{\partial r_{\kappa\alpha}^a}} u_{\kappa\alpha}^a + \frac{1}{2} \sum_{\substack{\kappa\alpha a \\ \kappa'\beta b}} \underbrace{C_{\kappa\alpha,\kappa'\beta}^{ab}}_{\frac{\partial^2 E}{\partial r_{\kappa\alpha}^a \partial r_{\kappa'\beta}^b}} u_{\kappa\alpha}^a u_{\kappa'\beta}^b$$



- Interatomic Force Constant (IFC) :
  - $\kappa \neq \kappa' : C_{\kappa\kappa'} = -K_{\kappa\kappa'}$
  - $\kappa = \kappa' : C_{\kappa\kappa} = -\sum' C_{\kappa\kappa'}$



Configuration space with  
 $3 \times N_{at}$  degrees of freedom

# Eigenmodes of vibration

- Equation of motion :

$$M_{\kappa} \ddot{u}_{\kappa\alpha}^a = - \sum_{\kappa'\beta b} C_{\kappa\alpha,\kappa'\beta}^{ab} u_{\kappa'\beta}^b$$

- Solution :

$$\begin{aligned} u_{\kappa\alpha}^a &= \eta_{\kappa\alpha}^a e^{-i\omega t} \\ &= \eta_{\kappa\alpha}^{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{R}_a} e^{-i\omega t} \end{aligned}$$

$$\omega^2 M_{\kappa} \eta_{\kappa\alpha}^{\mathbf{k}} = \sum_{\kappa'\beta} \eta_{\kappa'\beta}^{\mathbf{k}} \left( \sum_b C_{\kappa\alpha,\kappa'\beta}^{ab} e^{i\mathbf{k}\cdot(\mathbf{R}_b - \mathbf{R}_a)} \right)$$

$$\tilde{C}_{\kappa\alpha,\kappa'\beta}^{\mathbf{k}} = \sum_b C_{\kappa\alpha,\kappa'\beta}^{0b} e^{i\mathbf{k}\cdot\mathbf{R}_b}$$

$$\omega^2 M_{\kappa} \eta_{\kappa\alpha}^{\mathbf{k}} = \sum_{\kappa'\beta} \tilde{C}_{\kappa\alpha,\kappa'\beta}^{\mathbf{k}} \eta_{\kappa'\beta}^{\mathbf{k}}$$

“Generalized” eigenvalue problem

# Dynamical equation

$$\omega^2 M_{\kappa} \eta_{\kappa\alpha}^{\mathbf{k}} = \sum_{\kappa' \beta} \tilde{C}_{\kappa\alpha, \kappa' \beta}^{\mathbf{k}} \eta_{\kappa' \beta}^{\mathbf{k}}$$

- If we define :

$$\begin{aligned} \tilde{D}_{\kappa\alpha, \kappa' \beta}^{\mathbf{k}} &= \tilde{C}_{\kappa\alpha, \kappa' \beta}^{\mathbf{k}} / \sqrt{M_{\kappa} M_{\kappa'}} \\ \gamma_{\kappa\alpha}^{\mathbf{k}} &= \sqrt{M_{\kappa}} \eta_{\kappa\alpha}^{\mathbf{k}} \end{aligned}$$



$$\sum_{\kappa' \beta} \tilde{D}_{\kappa\alpha, \kappa' \beta}^{\mathbf{k}} \gamma_{\kappa' \beta}^{\mathbf{k}} = \omega^2 \gamma_{\kappa\alpha}^{\mathbf{k}}$$

*Dynamical equation*  
Classical eigenvalue problem

# Frequencies and eigendisplacements

ANADDB

$$\sum_{\kappa'\beta} \tilde{D}_{\kappa\alpha,\kappa'\beta}^{\mathbf{k}} \gamma_{\kappa'\beta}^{\mathbf{k}} = \omega^2 \gamma_{\kappa\alpha}^{\mathbf{k}}$$

Phonon frequency

Phonon eigenvector  $\langle \gamma^{\mathbf{k}} | \gamma^{\mathbf{k}} \rangle = 1$

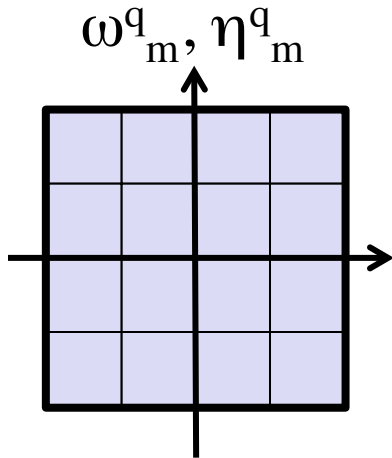
Phonon eigendisplacement  $\gamma_{\kappa\alpha}^{\mathbf{k}} = \sqrt{M_{\kappa}} \eta_{\kappa\alpha}^{\mathbf{k}}$   
 $\langle \eta^{\mathbf{k}} | M | \eta^{\mathbf{k}} \rangle = 1$

Dynamical matrix

$$\tilde{D}_{\kappa\alpha,\kappa'\beta}^{\mathbf{k}} = \frac{1}{\sqrt{M_{\kappa} M_{\kappa'}}} \frac{\partial^2 E}{\partial \eta_{\kappa\alpha}^{\mathbf{k}} \partial \eta_{\kappa'\beta}^{\mathbf{k}}}$$

Second energy derivatives from DFPT

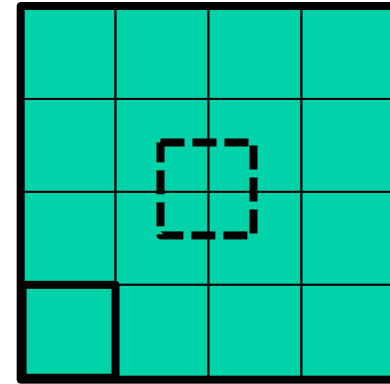
# From atom to phonon basis



m-bands :  $3x n_{\text{at}}$   
 q-points :  $N \times N \times N$  grid

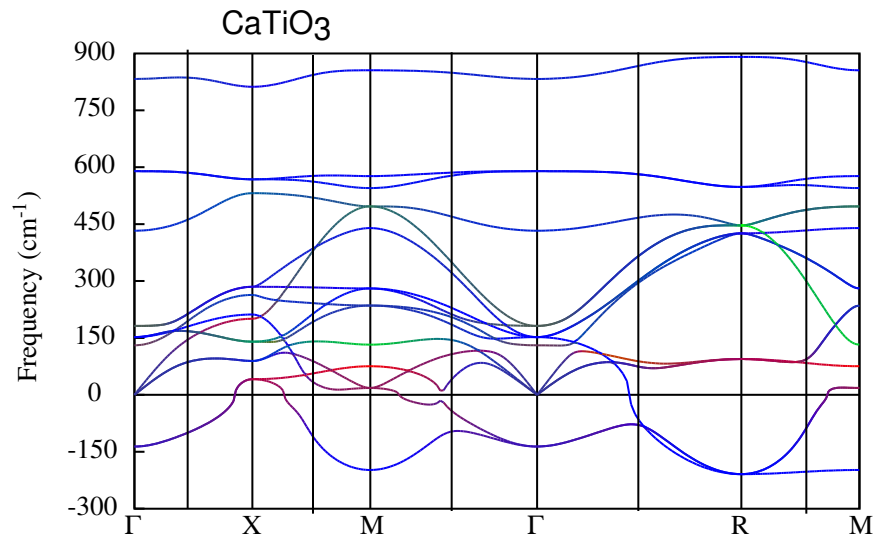
$$u_{\kappa\alpha}^a = \eta_{\kappa\alpha}^{\mathbf{k}} \cdot e^{i\mathbf{k} \cdot \mathbf{R}_a}$$

Commensurate  
with



$N \times N \times N$  supercell  
 $3x n_{\text{at}} \times N^3$  degrees of freedom

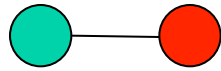
Phonon modes on the  $N^3$  grid are orthogonal to each others and form a complete basis for the atomic displacements in the  $N^3$  supercell





# From atom to phonon basis

**Atom basis**



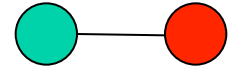
$$u_1 = (1, 0)$$



$$u_2 = (0, 1)$$



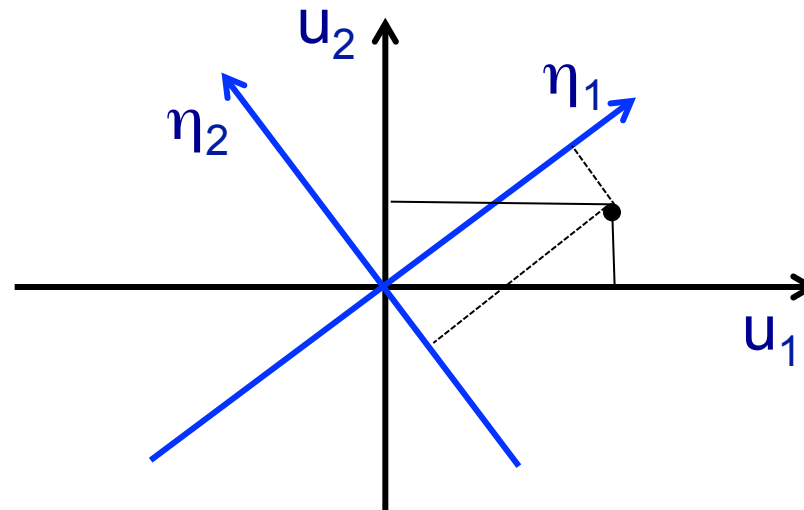
**Phonon basis**



$$\eta_1 = (1, 1) \sqrt{2}/2$$



$$\eta_2 = (-1, 1) \sqrt{2}/2$$

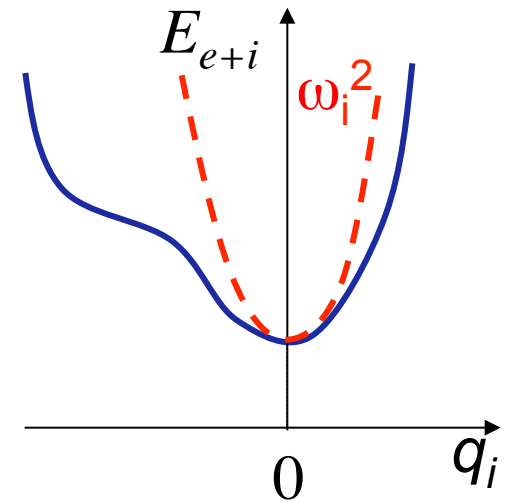


# From atom to

$3xN_{\text{at}}$  phonons form a complete basis :  $u_{\kappa\alpha}^a = \sum_i q_i \eta_{\kappa\alpha}^i$

$$\begin{aligned}
 E &= \frac{1}{2} \sum_{\substack{\kappa\alpha \\ \kappa'\beta}} C_{\kappa\alpha, \kappa'\beta} u_{\kappa\alpha} u_{\kappa'\beta} \\
 &= \frac{1}{2} \sum_{ij} q_i q_j \sum_{\substack{\kappa\alpha \\ \kappa'\beta}} C_{\kappa\alpha, \kappa'\beta} \eta_{\kappa\alpha}^i \eta_{\kappa'\beta}^j \\
 &= \frac{1}{2} \sum_{ij} q_i q_j \sum_{\kappa\alpha} \eta_{\kappa\alpha}^i \sum_{\kappa'\beta} C_{\kappa\alpha, \kappa'\beta} \eta_{\kappa'\beta}^j \\
 &= \frac{1}{2} \sum_{ij} q_i q_j \sum_{\kappa\alpha} \eta_{\kappa\alpha}^i M_{\kappa} \eta_{\kappa\alpha}^j \omega_j^2 \\
 &= \frac{1}{2} \sum_{ij} q_i q_j \omega_j^2 \delta_{ij} \\
 &= \frac{1}{2} \sum_i \omega_i^2 q_i^2
 \end{aligned}$$

Finite crystal



# Change of basis

$3xN_{\text{at}}$  phonons form a complete basis :  $u_{\kappa\alpha}^a = \sum_i q_i \eta_{\kappa\alpha}^i$

$$E_{e+i} = \frac{1}{2} \sum_{\substack{\kappa\alpha \\ \kappa'\beta}} C_{\kappa\alpha, \kappa'\beta} u_{\kappa\alpha} u_{\kappa'\beta}$$

$$= \frac{1}{2} \sum_{ij} q_i q_j \sum_{\substack{\kappa\alpha \\ \kappa'\beta}} C_{\kappa\alpha, \kappa'\beta} \eta_{\kappa\alpha}^i \eta_{\kappa'\beta}^j$$

$$= \frac{1}{2} \sum_{ij} q_i q_j$$

$$= \frac{1}{2} \sum_{ij} q_i q_j$$

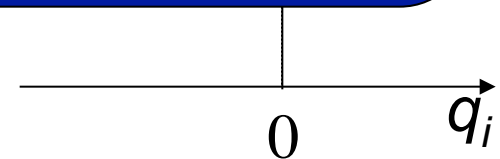
$$= \frac{1}{2} \sum_{ij} q_i q_j \omega_{ij}$$

$$= \frac{1}{2} \sum_i \omega_i^2 q_i^2$$

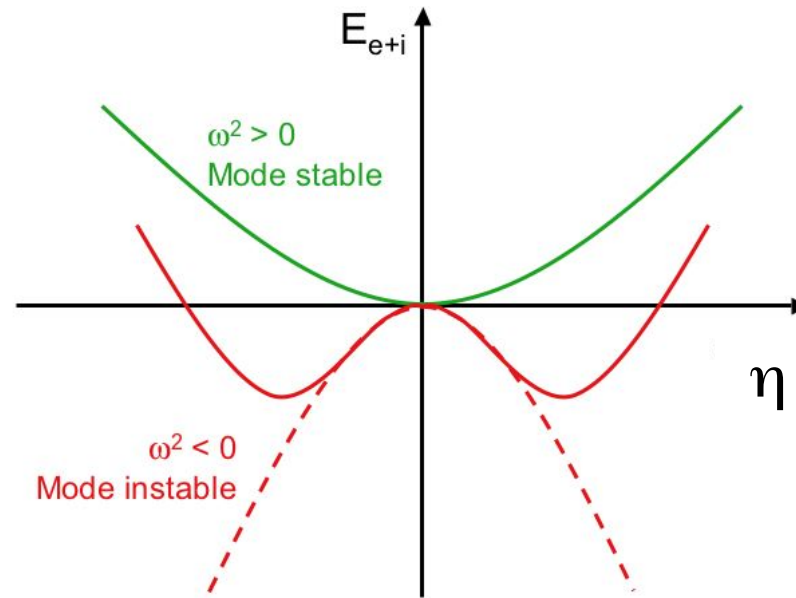
Finite crystal



1. Crystal = a collection of independent harmonic oscillators → phonons  
 2.  $\omega^2$  = curvature of  $E_{e+i}$  along a given phonon direction.



# Imaginary phonon frequencies



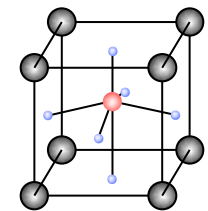
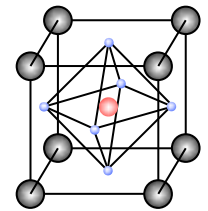
- Imaginary phonon frequencies ( $\omega^2 < 0$ ) are associated to structural instabilities.
- $\eta$  provides the direction to follow in order to reach the lower symmetry phase

Projection on phonon modes

# Motivation

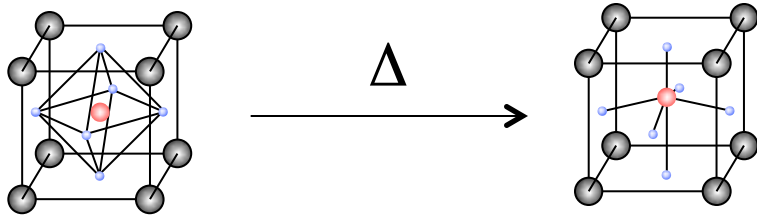
- In crystals showing a displacive phase transition (like mode  $ABO_3$  perovskites), the low symmetry phase appears as a small distortion of a (eventually hypothetical) high-symmetry reference.
- It is possible to analyse the atomic distortion in the low symmetry phase by projecting on the phonon of the high-symmetry phase and get many insight on the phase transition

Cubic  
reference



Distorted  
phase

# Mode-by-mode decomposition



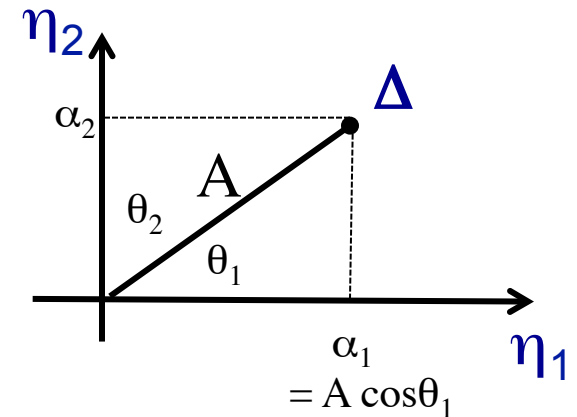
**Compute distortion vector  $\Delta$  :**

- in reduced coordinates
- fixing CM

**Mode contribution :**

$$\Delta = \sum_i \alpha_i \eta_i = A \sum_i \cos \theta_i \eta_i$$

$$\langle \eta_j | M | \Delta \rangle = \sum_i \alpha_i \langle \eta_j | M | \eta_i \rangle = \alpha_j = A \cos \theta_j$$



**Typically we can report :**

- The amplitude of each mode:  $\alpha_i$
- The distortion amplitude : A
- The relative mode contribution :  $\cos \theta_i$

$$\begin{aligned} \langle \Delta | M | \Delta \rangle &= \sum_i \alpha_i^2 = A^2 \\ \sum_i \alpha_i^2 &= A^2 \\ \sum_i (\cos \theta_i)^2 &= 1 \end{aligned}$$

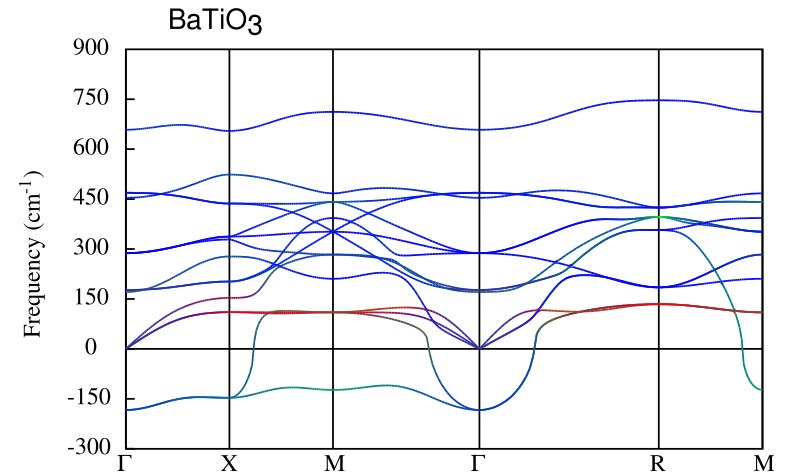
# Ferroelectric distortion BaTiO<sub>3</sub>

Paraelectric

Ferroelectric



	<i>P4mm</i>	<i>Amm2</i>	<i>R3m</i>
BaTiO <sub>3</sub>			
$ \boldsymbol{\tau} $	0.188	0.209	0.214
$F_{1u}(\text{TO}_1)$	0.993	0.975	0.971
$F_{1u}(\text{TO}_2)$	0.110	0.087	0.129
$F_{1u}(\text{TO}_3)$	0.032	0.194	0.228
CaTiO <sub>3</sub>			
$ \boldsymbol{\tau} $	0.601	0.478	0.435
$F_{1u}(\text{TO}_1)$	0.985	0.970	0.970
$F_{1u}(\text{TO}_2)$	-0.171	-0.203	-0.199
$F_{1u}(\text{TO}_3)$	0.033	0.129	-0.143
CaZrO <sub>3</sub>			
$ \boldsymbol{\tau} $	0.903	0.886	0.868
$F_{1u}(\text{TO}_1)$	0.976	0.910	0.897
$F_{1u}(\text{TO}_2)$	-0.060	0.202	-0.208
$F_{1u}(\text{TO}_3)$	-0.067	-0.336	-0.197

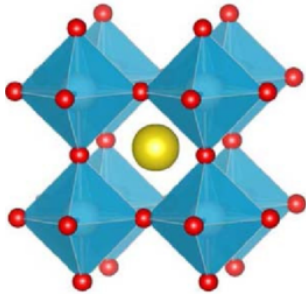


Strong domination of the unstable mode

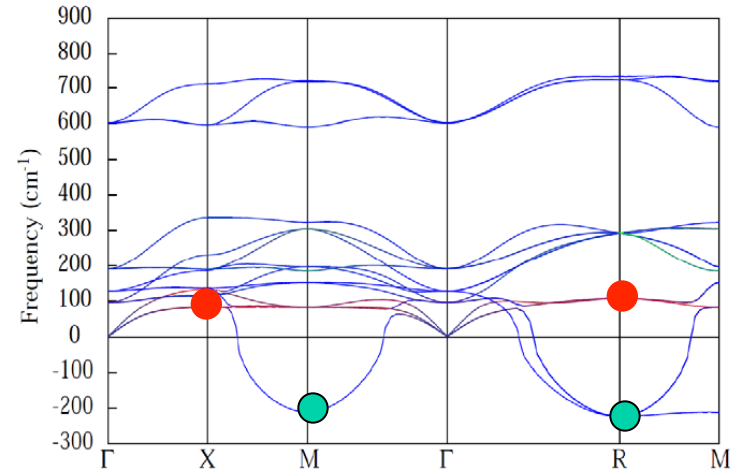
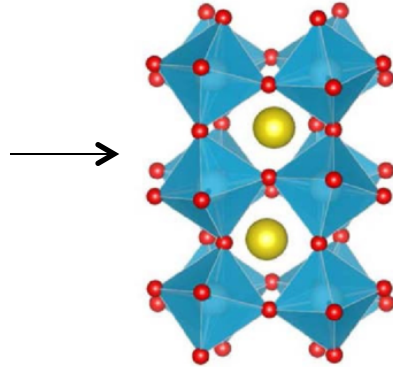


# Antiferrodistortive distortion SrRuO<sub>3</sub>

Cubic



Orthorhombic



Rotations	Symmetry	Space group number	Oxygen rotations		Anti-polar motions		Other oxygen motions	
			R <sub>4</sub> <sup>+</sup>	M <sub>3</sub> <sup>+</sup>	X <sub>5</sub> <sup>+</sup>	R <sub>5</sub> <sup>+</sup>	M <sub>2</sub> <sup>+</sup>	M <sub>4</sub> <sup>+</sup>
<i>a<sup>0</sup>a<sup>0</sup>c<sup>-</sup></i>	<i>I4/mcm</i>	140	0.73	—	—	—	—	—
<i>a<sup>0</sup>a<sup>0</sup>c<sup>+</sup></i>	<i>P4/mbm</i>	127	—	0.71	—	—	—	—
<i>a<sup>0</sup>b<sup>-</sup>b<sup>-</sup></i>	<i>Imma</i>	74	0.80	—	—	0.08	—	—
<i>a<sup>0</sup>b<sup>+</sup>b<sup>+</sup></i>	<i>I4/mmm</i>	139	—	1.05	—	—	—	0.09
<i>a<sup>0</sup>b<sup>+</sup>c<sup>-</sup></i>	<i>Cmcm</i>	63	0.85	0.81	0.20	0.03	—	0.04
<i>a<sup>-</sup>b<sup>+</sup>a<sup>-</sup></i>	<i>Pnma</i>	62	1.00	0.77	0.35	0.06	0.00	—
<i>a<sup>+</sup>a<sup>+</sup>c<sup>-</sup></i>	<i>P4(2)/nmc</i>	137	1.15	1.18	0.30	—	—	0.07

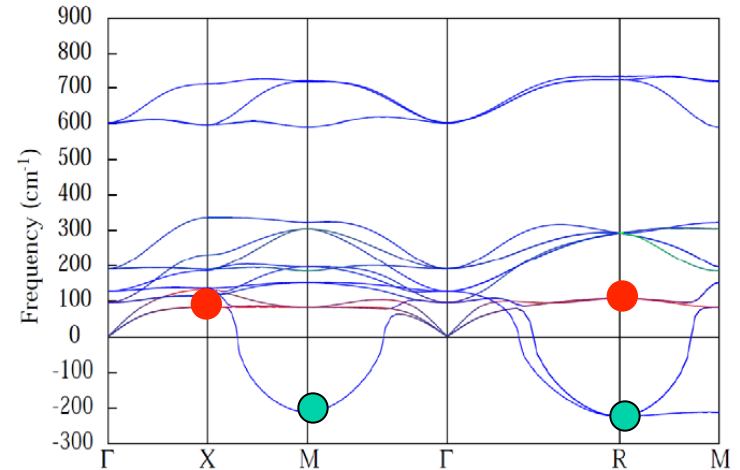
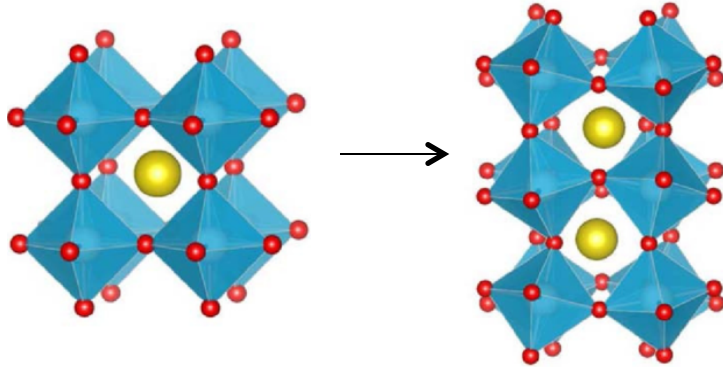
The unstable modes are dominant BUT other stable modes appear due to anharmonic couplings !

$$a^-b^+a^- \quad Pnma \quad 62 \quad Q(R_4^+)Q(M_3^+)Q(X_5^+) + Q(R_4^+)Q(M_2^+)Q(X_5^+) + Q(R_5^+)Q(M_3^+)Q(X_5^+) \dots$$

# Antiferrodistortive distortion SrRuO<sub>3</sub>

Cubic

Orthorhombic



Rotations	Symmetry	Space group number
$a^0a^0c^-$	$I4/mcm$	140
$a^0a^0c^+$	$P4/mbm$	127
$a^0b^-b^-$	$Imma$	74
$a^0b^+b^+$	$I4/mmm$	139
$a^0b^+c^-$	$Cmcm$	63
$a^-b^+a^-$	$Pnma$	62
$a^+a^+c^-$	$P4(2)/nmc$	137

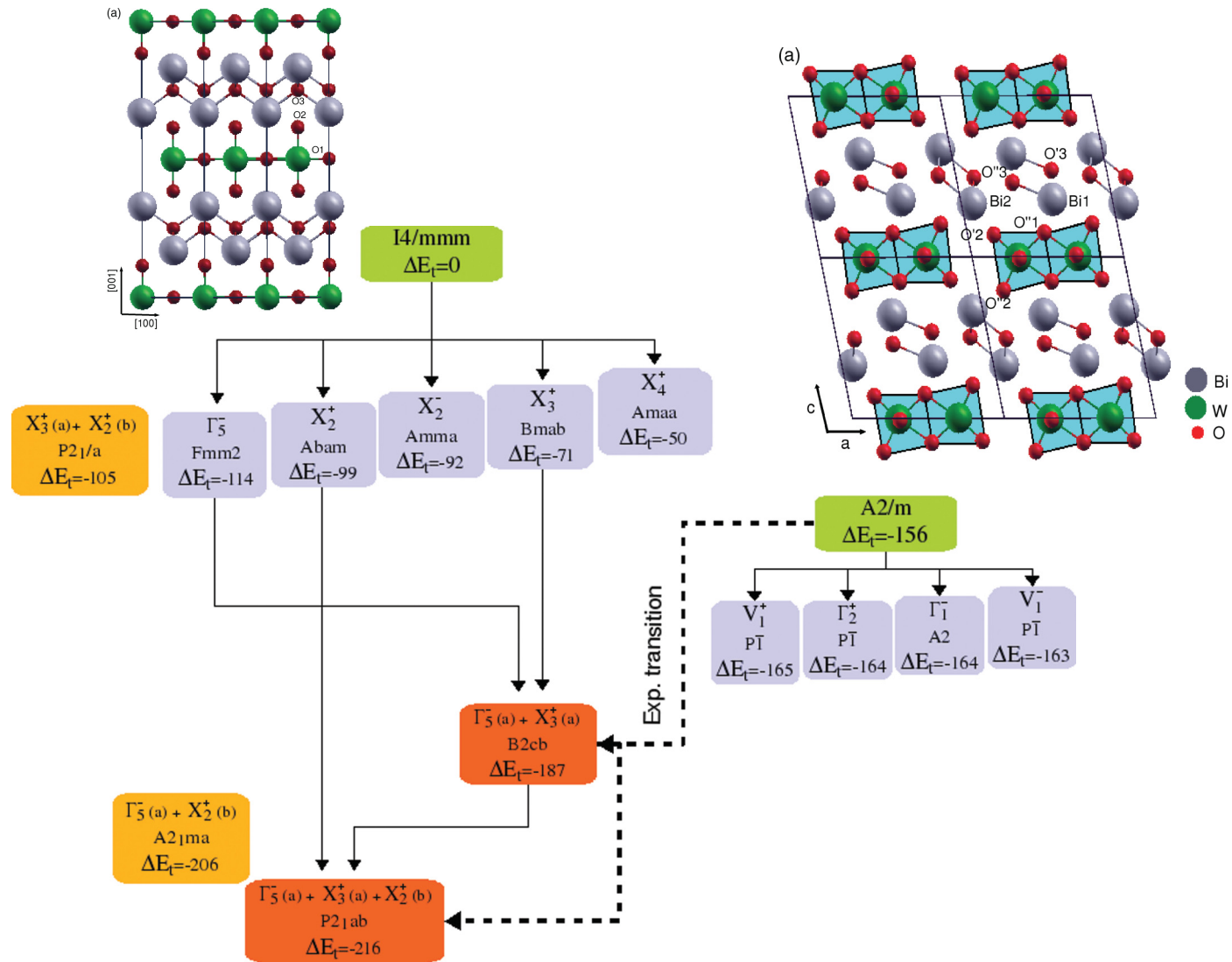
Oxygen rotations    Anti-polar motions    Other oxygen motions

$$\begin{aligned}
 E = & \frac{1}{2} A_1 Q_M^2 + \frac{1}{4} B_1 Q_M^4 \\
 & + \frac{1}{2} A_2 Q_R^2 + \frac{1}{4} B_2 Q_R^4 + C_{12} Q_R^2 Q_M^2 \\
 & + \frac{1}{2} A_3 Q_X^2 - \lambda Q_M Q_R Q_X + \dots
 \end{aligned}$$

The unstable modes are dorn  
stable modes appear due to

$$a^-b^+a^- \quad Pnma \quad 62 \quad Q(R_4^+)Q(M_3^+)Q(X_5^+) + Q(R_4^+)Q(M_2^+)Q(X_5^+) + Q(R_5^+)Q(M_3^+)Q(X_5^+) \dots$$

# Aurivillius $\text{Bi}_2\text{WO}_6$



# Aurivillius $\text{Bi}_2\text{WO}_6$

Phase	A	$\Gamma_5^-$ [198i]	$X_2^-$ [183i]	$X_2^+$ [135i]	$X_3^+$ [104i]	$X_4^+$ [98i]	$\Gamma_5^-$ [29]	$X_2^-$ [80]	$X_3^+$ [107]	$\Gamma_5^-$ [137]
Phases arising from single-mode condensation										
<i>Fmm2</i>	361.0	0.84					0.51			0.11
<i>Amma</i>	395.6		0.79					0.60		
<i>Abam</i>	444.7			0.99						
<i>Bmab</i>	444.5				0.98				0.14	
<i>Amaa</i>	384.0					0.99				
Experimentally observed phases										
<i>B2cb</i>	601.0	0.53			0.79		0.24			0.09
<i>P2<sub>1</sub>ab</i>	582.7	0.52		0.59	0.55		0.04		0.1	0.09
Hypothetical phases										
<i>A2<sub>1</sub>ma</i>	514.5	0.58		0.78			0.01			0.11
<i>P2<sub>1</sub>/a</i>	455.0			0.76	0.62				0.11	

Allow to follow mode amplitude evolution from one phase to another and to understand mode competition or cooperation

$$\begin{aligned}
 \Delta E_t(Q_{\Gamma_5^-}, Q_{X_2^+}, Q_{X_3^+}) &= -2.50Q_{\Gamma_5^-}^2 - 7.50 \times 10^{-1}Q_{X_3^+}^2 - 1.03Q_{X_2^+}^2 \\
 &+ 1.36 \times 10^{-5}Q_{\Gamma_5^-}^4 + 1.97 \times 10^{-6}Q_{X_3^+}^4 \\
 &+ 2.65 \times 10^{-6}Q_{X_2^+}^4 + 2.10 \times 10^{-6}Q_{\Gamma_5^-}^2Q_{X_3^+}^2 \\
 &- 1.79 \times 10^{-9}Q_{\Gamma_5^-}^2Q_{X_2^+}^2 + 1.79 \times 10^{-6}Q_{X_3^+}^2Q_{X_2^+}^2.
 \end{aligned}$$

# Aurivillius $\text{Bi}_2\text{WO}_6$

## Comparison of theoretical and experimental GS structures

### Atomic positions

Structure	Atom	Wyckoff	$x$	$y$	$z$
$P2_1ab$					
$a = 5.30 \text{ \AA}$ (5.45)	$\text{Bi}_1$	$4a$	0.0109	0.4843	0.1698
			(-0.0126	0.5113	0.1730)
$b = 5.32 \text{ \AA}$ (5.48)	$\text{Bi}_2$	$4a$	0.0036	0.5074	-0.1698
			(-0.0113	0.4762	0.1719)
$c = 16.17 \text{ \AA}$ (16.47)	W	$4a$	0.0000	0.0000	0.0000
			$\text{O}'_1$	$4a$	0.2318
	$\text{O}''_1$	$4a$	(0.2679	0.6933	-0.0147)
			0.3321	0.1984	-0.0109
	$\text{O}'_2$	$4a$	(0.3342	0.2220	0.0163)
			-0.0684	-0.0368	0.1109
	$\text{O}''_2$	$4a$	(0.0854	0.0449	0.1072)
			0.4428	0.4563	-0.1108
	$\text{O}'_3$	$4a$	(0.5703	0.5526	-0.1086)
			0.2676	0.2413	0.7515
$\text{O}''_3$	$4a$	(0.2728	0.2508	0.7489)	
		0.2752	0.2465	0.2516	
			(0.2740	0.2326	0.2514)

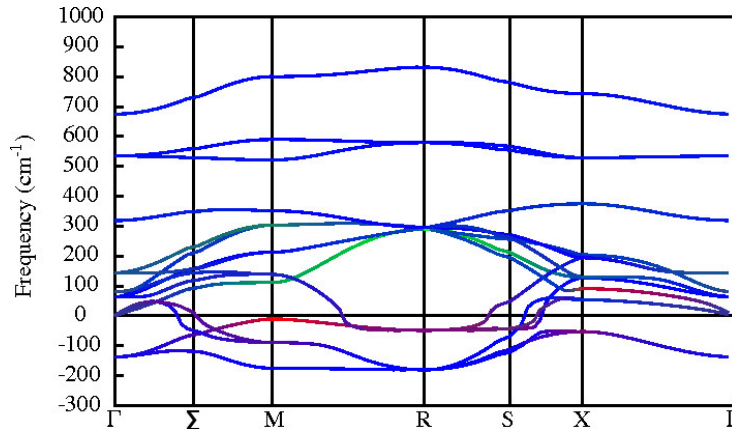
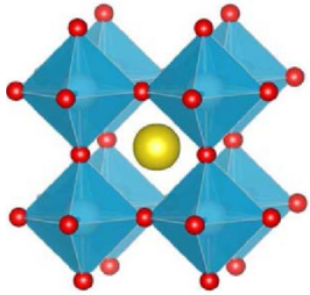
### Mode décomposition

**Table 2. Contribution of the Phonon Modes of the Hypothetical  $I4/mmm$  Phase (Frequencies in Brackets in  $\text{cm}^{-1}$ ) to the Distortion from the  $I4/mmm$  Phase to the Calculated and Experimental (300 K)  $P2_1ab$  Ground State<sup>a</sup>**

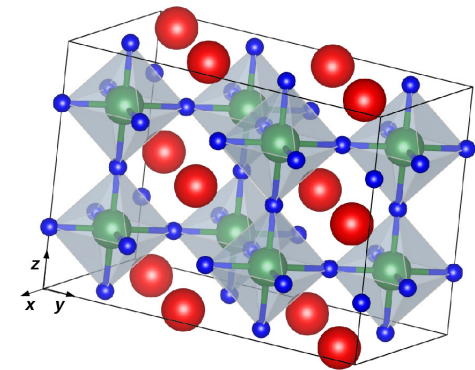
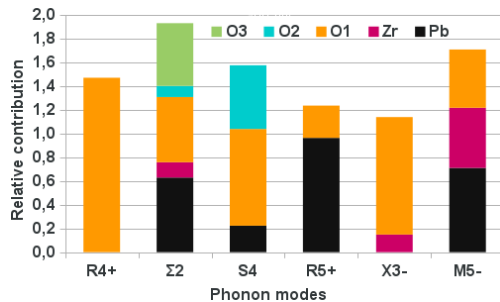
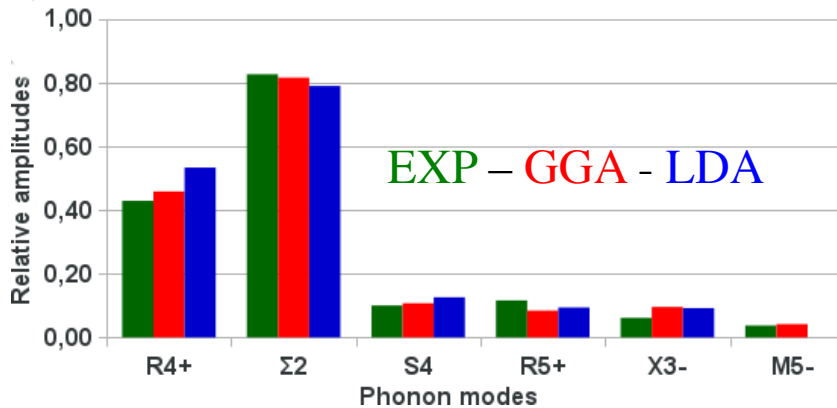
	A	$\Gamma_5^-$ (198i)	$X_2^+$ (135i)	$X_3^+$ (104i)	$\Gamma_5^-$ (29)
calculated (present)	585	0.52	0.59	0.55	0.04
experiment (300 K) <sup>6</sup>	772	0.56	0.30	0.59	0.35

<sup>a</sup>Only the most significant contributions are retained.  $A$  is the distortion amplitude (in Bohr) and  $\alpha_i$  the cosine director, according to the conventions and normalization conditions defined in ref 7.

# Antiferroelectric PbZrO<sub>3</sub>



<i>k</i> -points	Mode	Type	$\omega$	
			GGA-WC	LDA
$\Gamma$ (0, 0, 0)	$\Gamma_4^-$	(FE)	137 <i>i</i>	143 <i>i</i>
<i>X</i> ( $\frac{1}{2}$ , 0, 0)	$X_5^+$	(AP)	52 <i>i</i>	58 <i>i</i>
<i>M</i> ( $\frac{1}{2}$ , $\frac{1}{2}$ , 0)	$M_3^+$	(AFD)	175 <i>i</i>	190 <i>i</i>
	$M_5^-$	(AP)	89 <i>i</i>	100 <i>i</i>
	$M_2^-$	(AP)	10 <i>i</i>	16 <i>i</i>
<i>R</i> ( $\frac{1}{2}$ , $\frac{1}{2}$ , $\frac{1}{2}$ )	$R_4^+$	(AFD)	180 <i>i</i>	196 <i>i</i>
	$R_5^+$	(AP)	48 <i>i</i>	53 <i>i</i>
$\Sigma$ ( $\frac{1}{4}$ , $\frac{1}{4}$ , 0)	$\Sigma_2$	(AP)	121 <i>i</i>	
	$\Sigma_2$	(AP)	62 <i>i</i>	
<i>S</i> ( $\frac{1}{2}$ , $\frac{1}{4}$ , $\frac{1}{4}$ )	$S_4$	(AFD/AP)	128 <i>i</i>	
	$S_4$	(AFD/AP)	50 <i>i</i>	



# qAgate

(**A**binet **G**raphical **A**nalysis **T**echnical **E**ngine)

## Amongst many other functionalities :

- Condense phonon modes with given amplitudes in a REF structure (useful to look at energy wells or fit an energy expansion in terms of selected degrees of freedom).
- Provides mode-by-mode analysis of a the distortion between a REF and a distorted structure.

(a public version with documentation soon)